

Simulations of the Ising model on a Shastry-Sutherland lattice by quantum annealing

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ENERGY

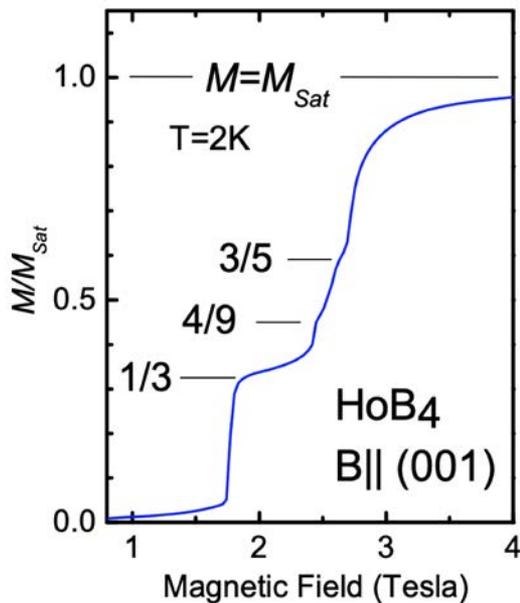
Outline

- Material science motivation
- Shastry-Sutherland lattice
- Embeddings
- Boundary conditions
- Information processing on D-Wave 2000Q
- Results and comparison to experimental data

We demonstrate how quantum annealing enables accurate simulations of many-particle Hamiltonian systems.

Rare-earth tetraboride materials

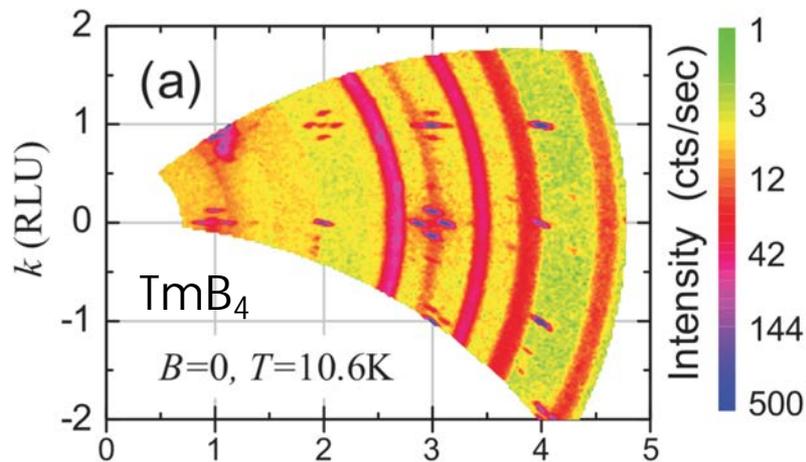
Magnetization



S Mat'as'et al. 2010

Structure Factor

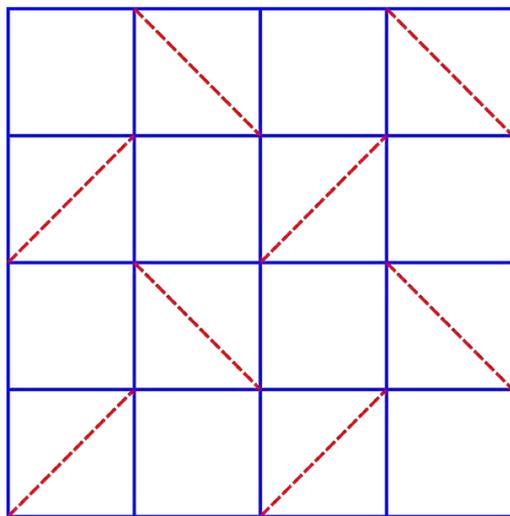
$$S(\vec{q}) = \sum_{\langle ij \rangle} \langle \sigma_i^z \sigma_j^z \rangle e^{i\vec{q} \cdot \vec{R}_{ij}}$$



Siemensmeyer et al. 2008

Can we simulate these behaviors using quantum annealing?

The Shastry-Sutherland lattice



Ising magnet

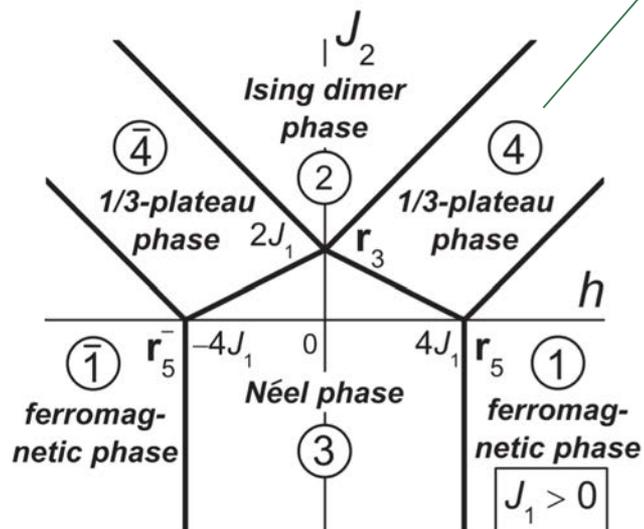
$$H = J_1 \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z + J_2 \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z$$

\longleftrightarrow RB₄, R = La-Lu

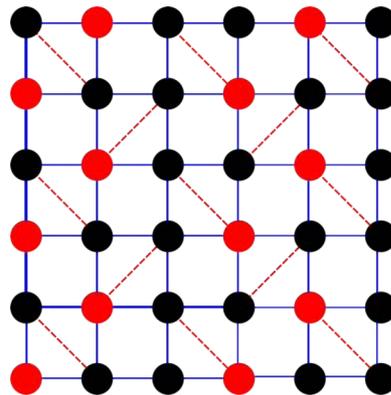
$J_2 = \text{Dimer}$

$J_1 = \text{Square}$

Phase diagram

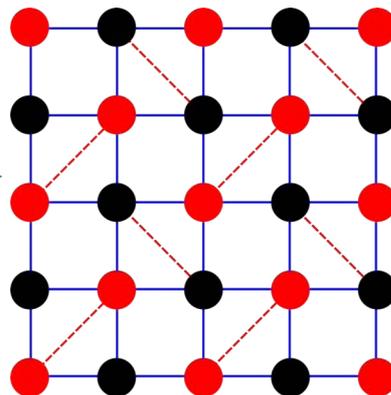
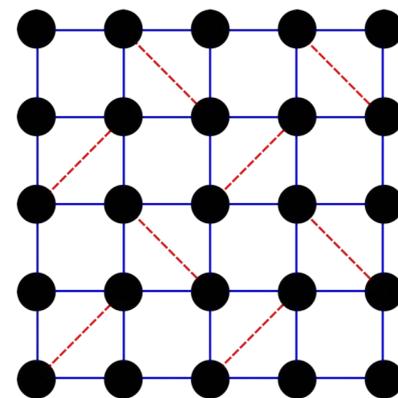


Dublenych 2012



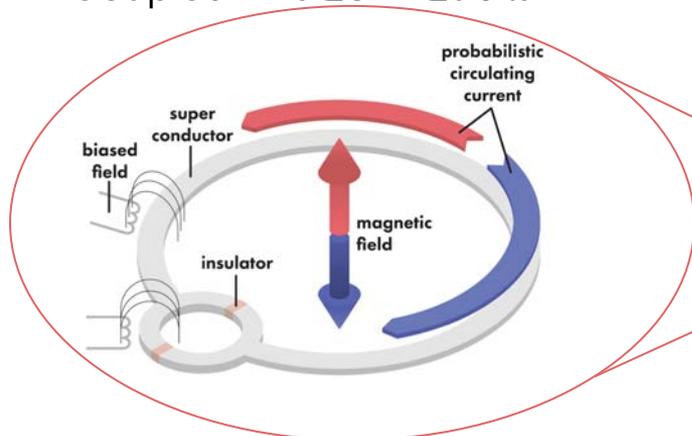
● Up spin

● Down spin



D-Wave 2000Q Processor

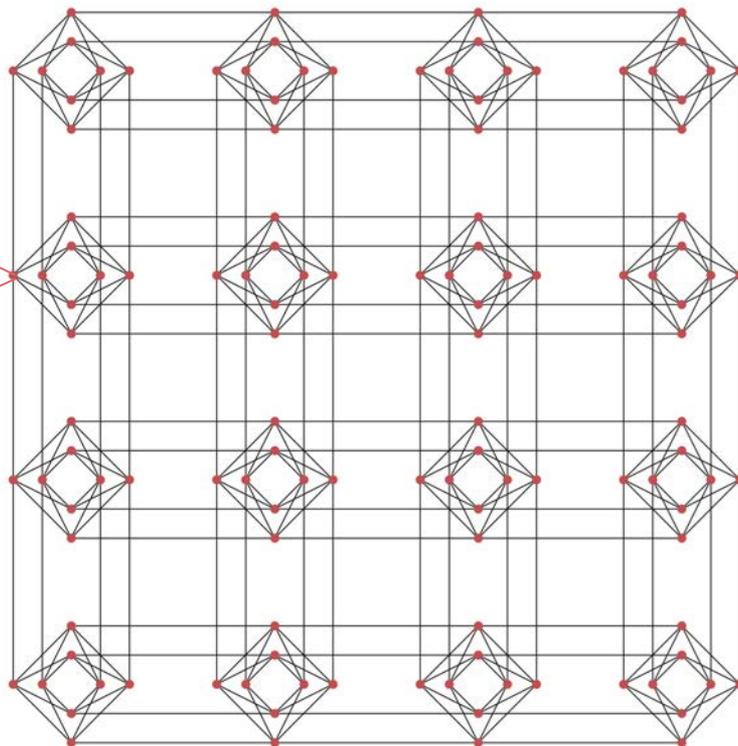
Coupled RF-SQUID Qubits



$$H = A(s) \sum_i \sigma_i^x + B(s) \left[\sum_i h_i \sigma_i^z + \sum_{\langle ij \rangle} J_{ij} \sigma_i^z \sigma_j^z \right]$$

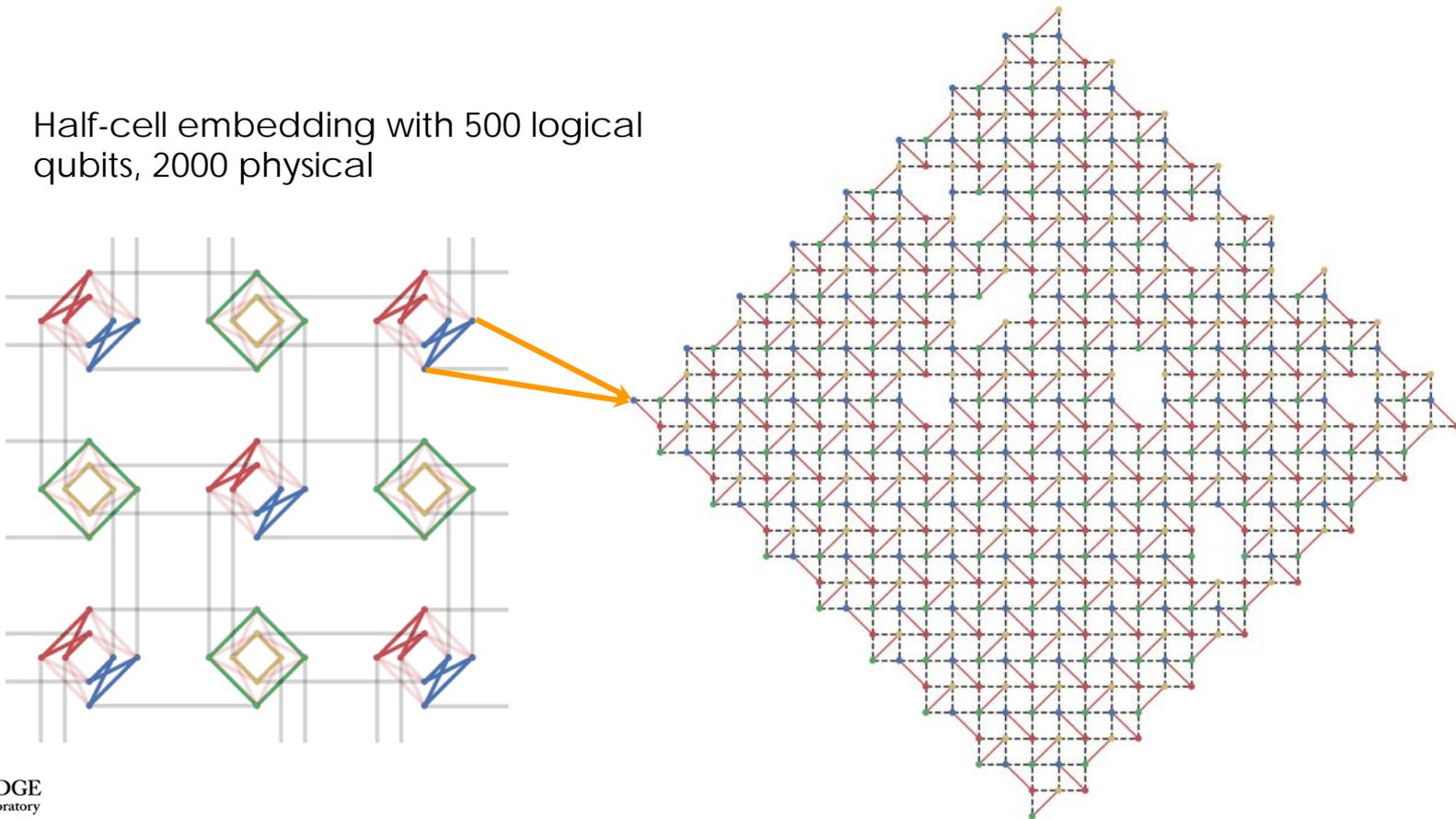
Independently tunable biases and couplers

Chimera Topology



Embedding the Shastry-Sutherland Lattice

- Half-cell embedding with 500 logical qubits, 2000 physical

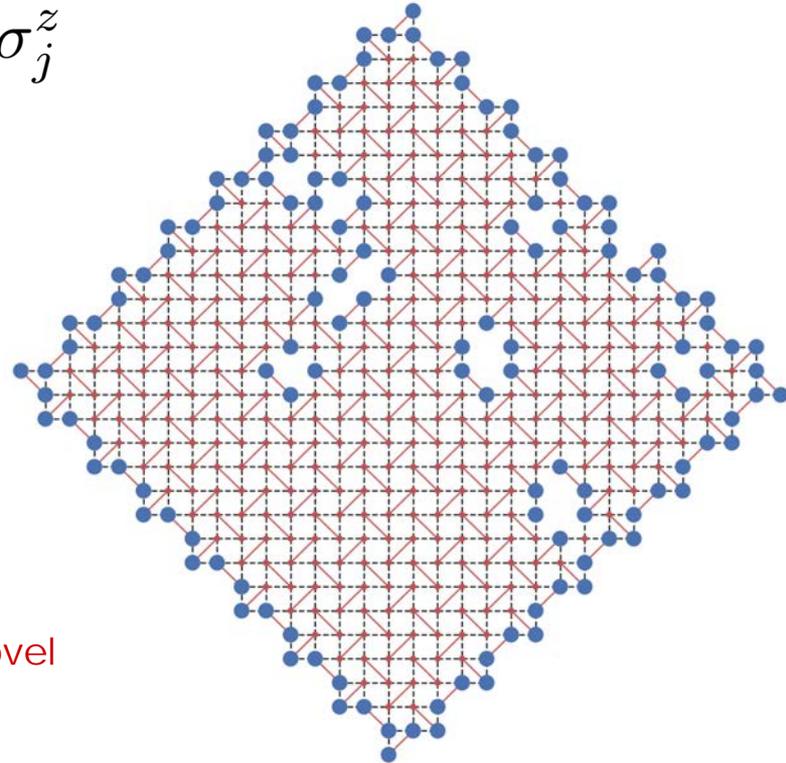


Self-consistent mean-field boundary conditions

$$H = \sum_i h_i^z \sigma_i^z + h^z \sum_i \sigma_i^z + \sum_{\langle ij \rangle} J_{ij} \sigma_i^z \sigma_j^z$$

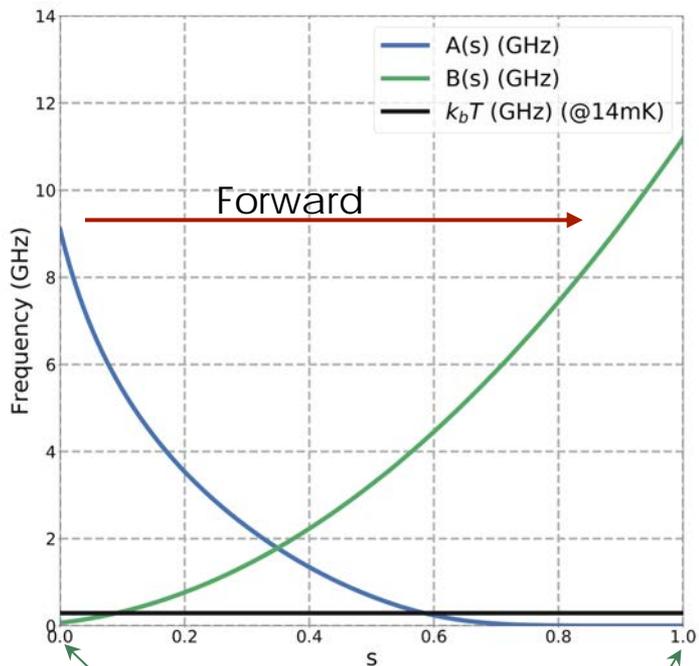
minimize $\langle m \rangle - \langle m \rangle$
 \vec{h}

subject to $\text{sgn}(h_i^z) = \text{sgn}(h^z) \quad \forall i$



Independent tunability of biases allows us to explore novel boundary conditions

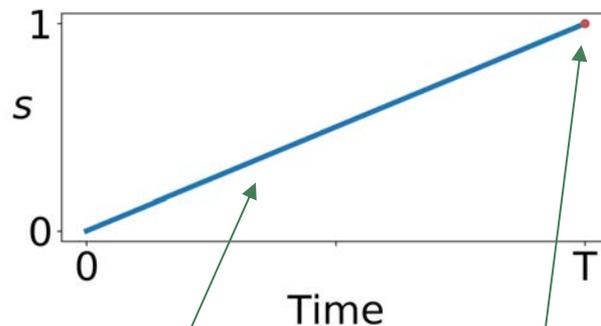
Forward Annealing



Begin
 $A(s) \gg B(s)$

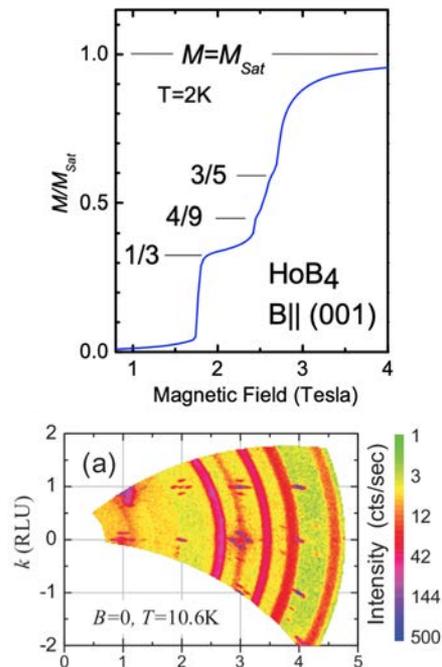
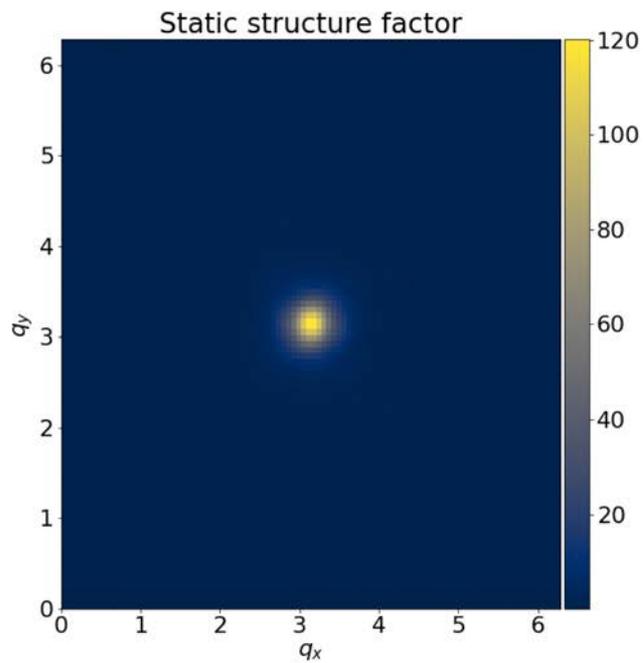
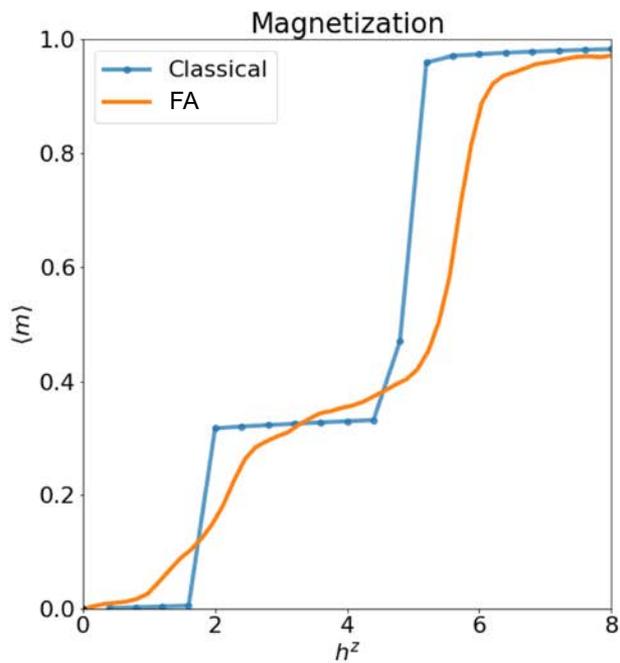
End
 $A(s) \ll B(s)$

$$H = A(s) \sum_i \sigma_i^x + B(s) \left[\sum_i h_i \sigma_i^z + \sum_{\langle ij \rangle} J_{ij} \sigma_i^z \sigma_j^z \right]$$

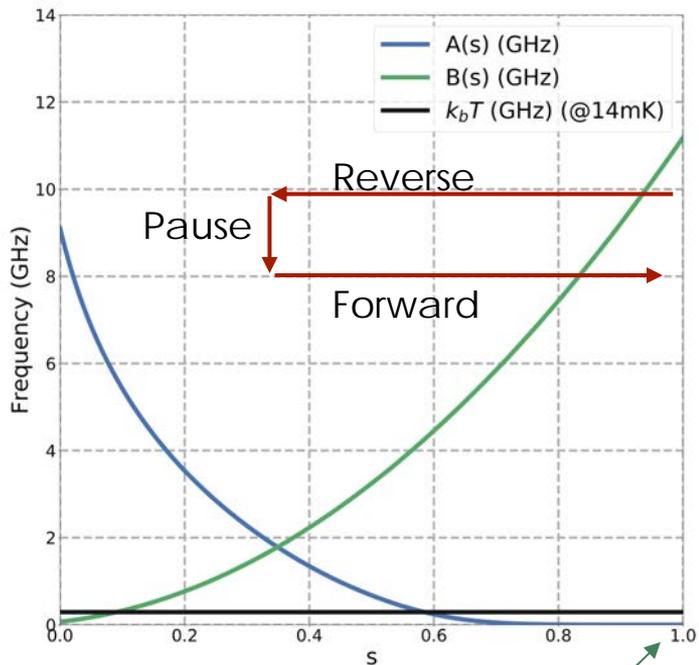


Adiabatic evolution and readout

Forward Anneal with the half-cell embedding

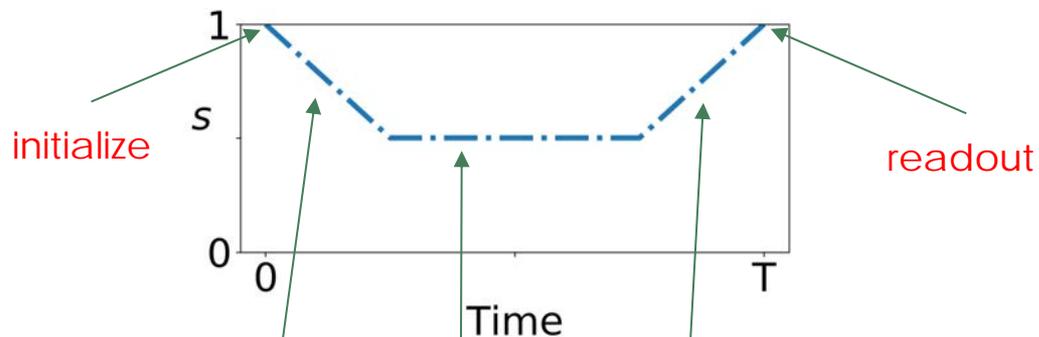


Reverse Annealing



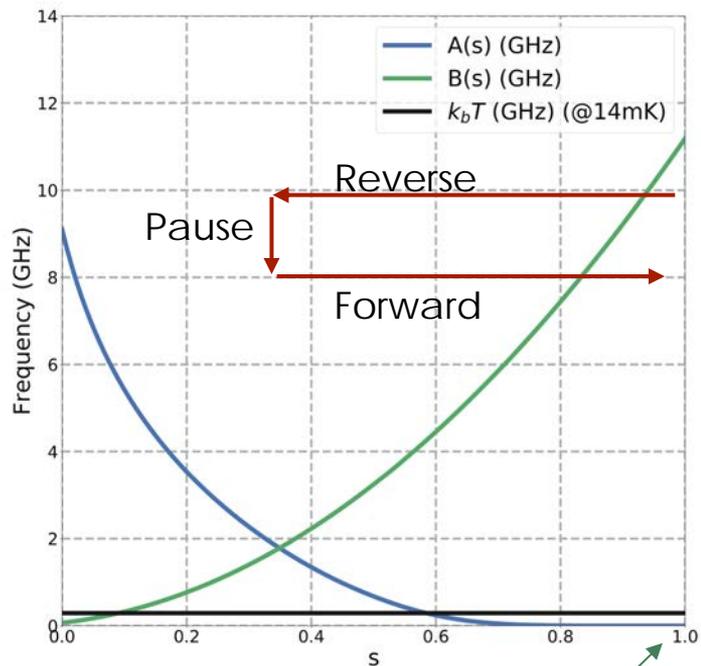
Begin + End
 $A(s) \ll B(s)$

$$H = A(s) \sum_i \sigma_i^x + B(s) \left[\sum_i h_i \sigma_i^z + \sum_{\langle ij \rangle} J_{ij} \sigma_i^z \sigma_j^z \right]$$



Reverse evolution , pause, forward evolution

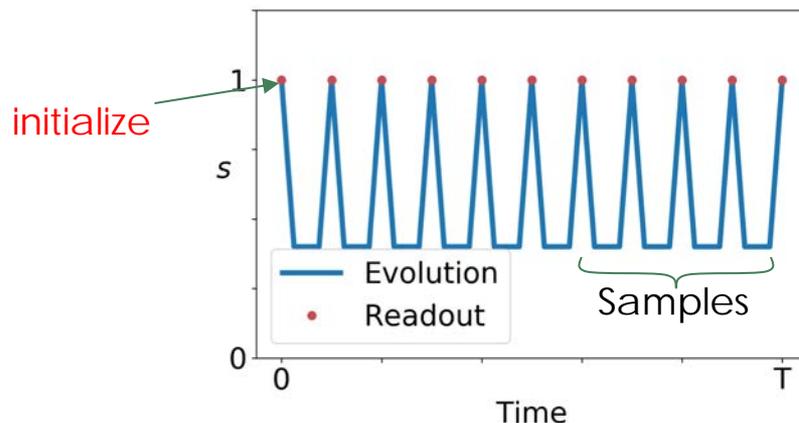
Quantum Evolution Markov Chain



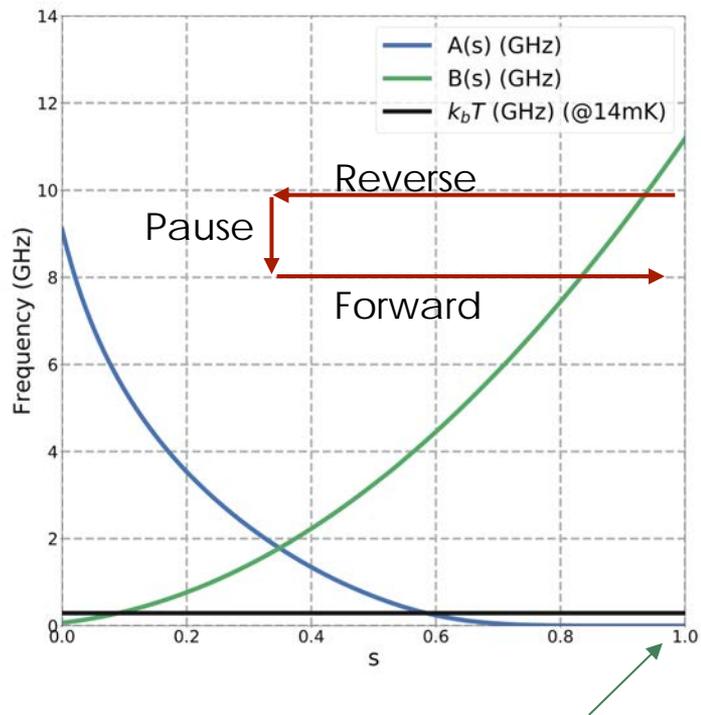
Begin + End
 $A(s) \ll B(s)$

$$H = A(s) \sum_i \sigma_i^x + B(s) \left[\sum_i h_i \sigma_i^z + \sum_{\langle ij \rangle} J_{ij} \sigma_i^z \sigma_j^z \right]$$

- Iterative reverse annealing schedule



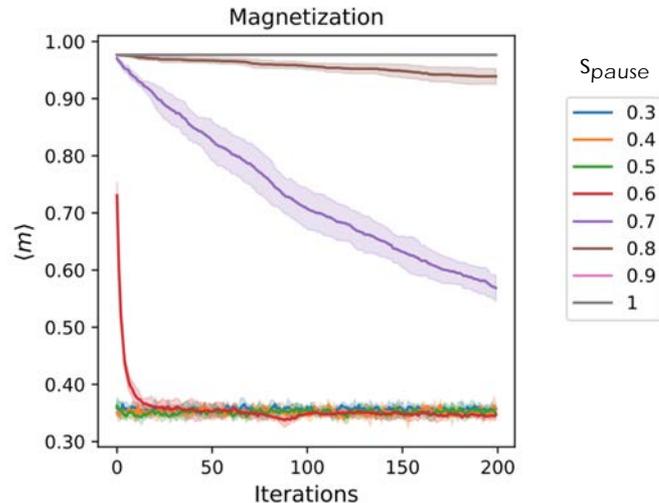
Quantum evolution Markov chain



Begin + End
 $A(s) \ll B(s)$

$$H = A(s) \sum_i \sigma_i^x + B(s) \left[\sum_i h_i \sigma_i^z + \sum_{\langle ij \rangle} J_{ij} \sigma_i^z \sigma_j^z \right]$$

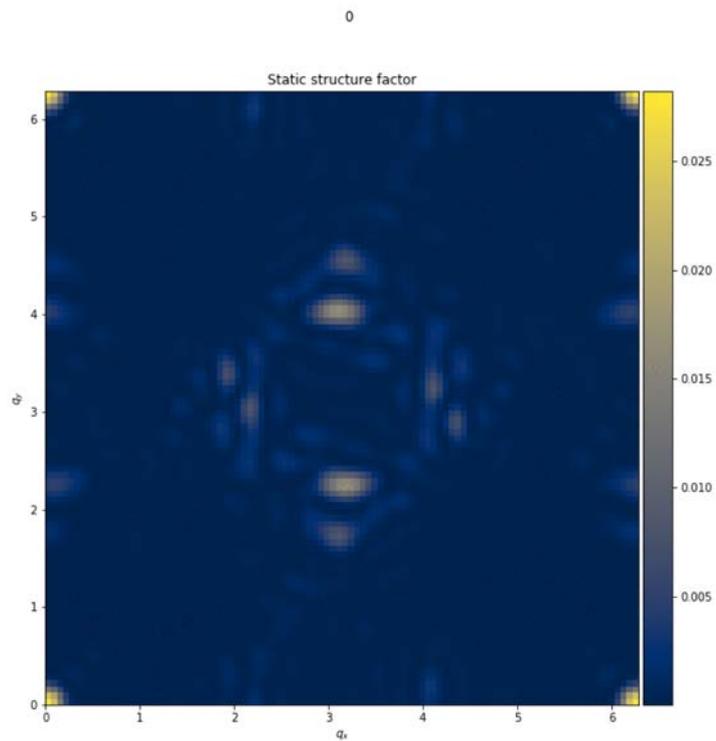
- Iterative reverse annealing schedule



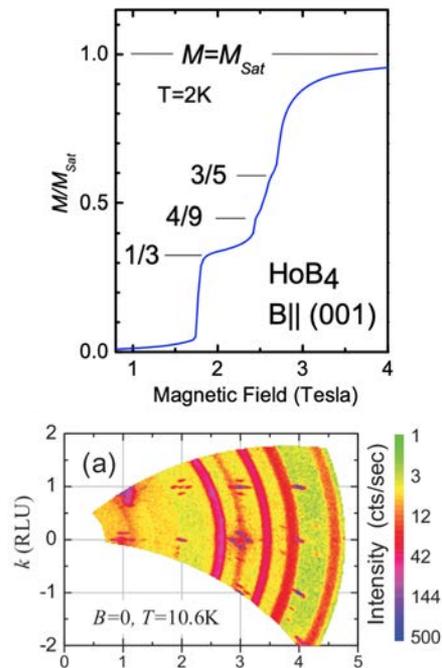
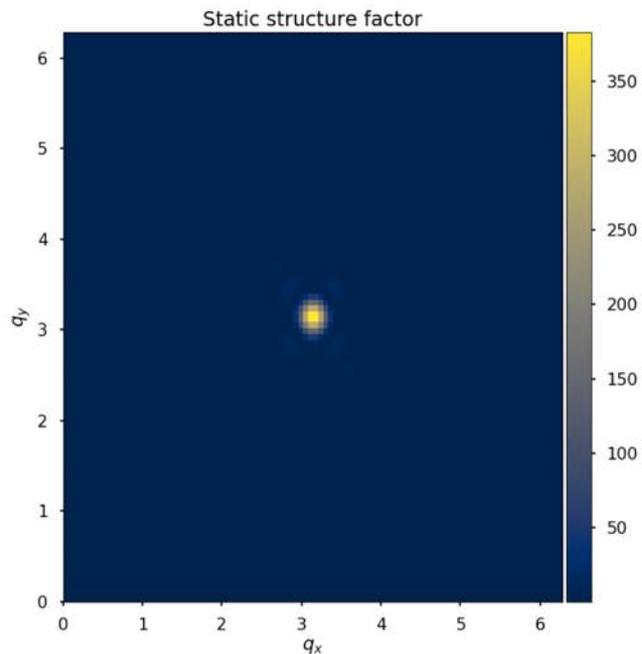
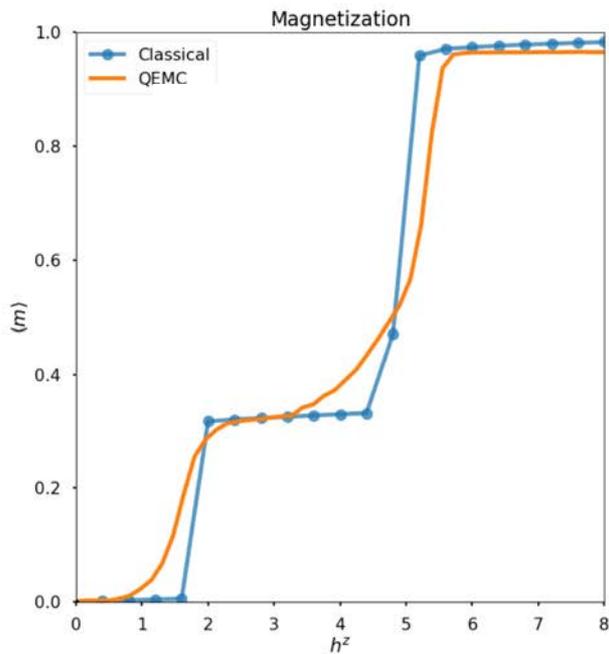
QEMC Motif Convergence

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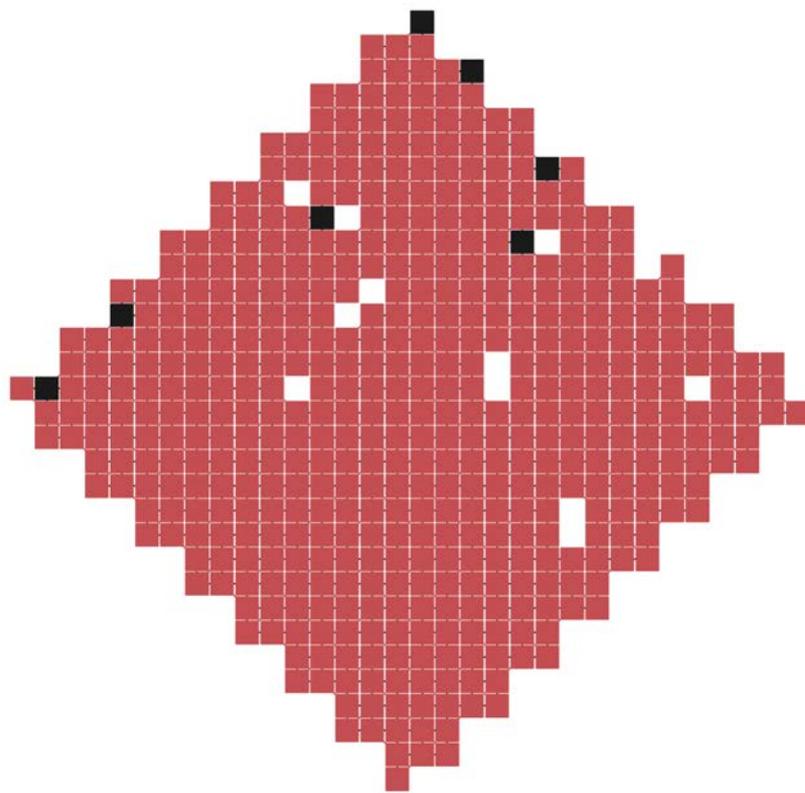
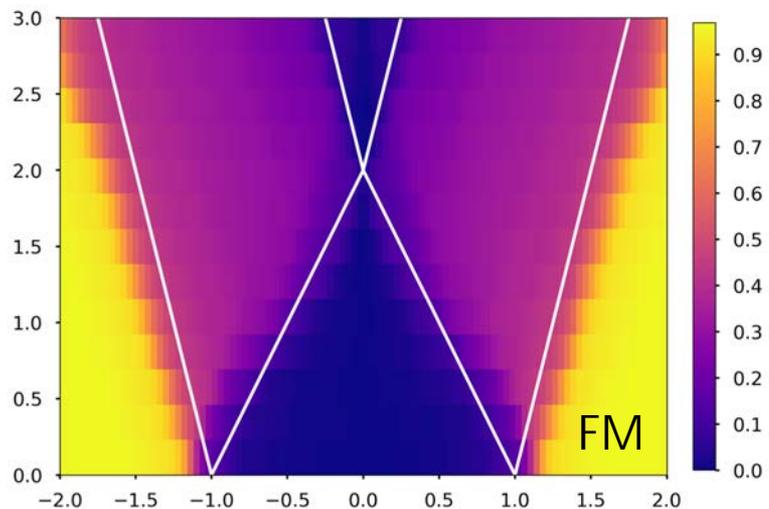
QEMC Structure Factor convergence



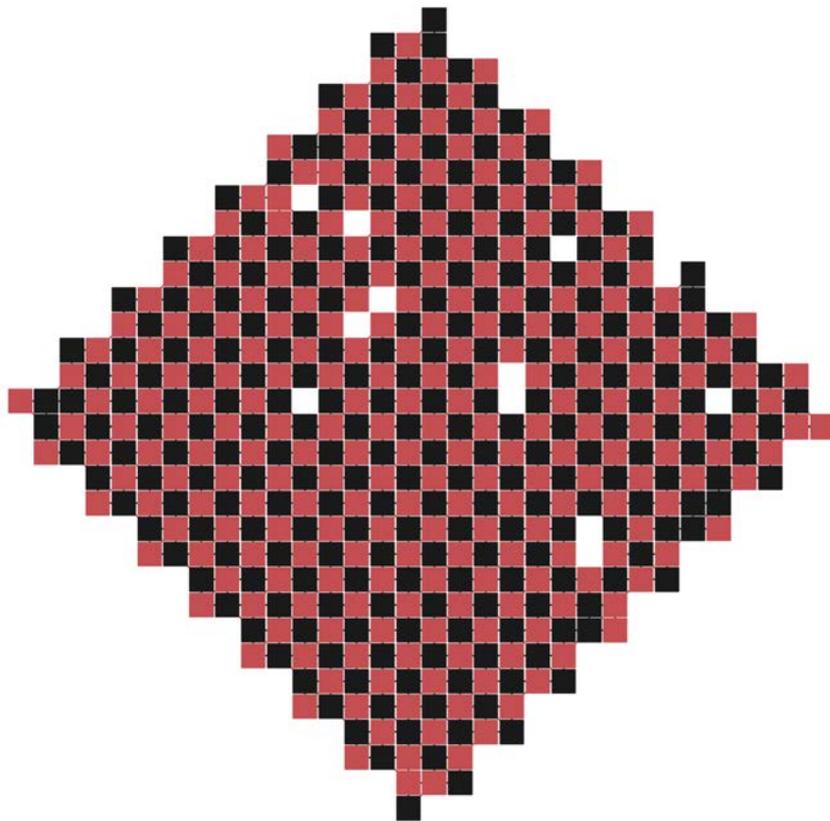
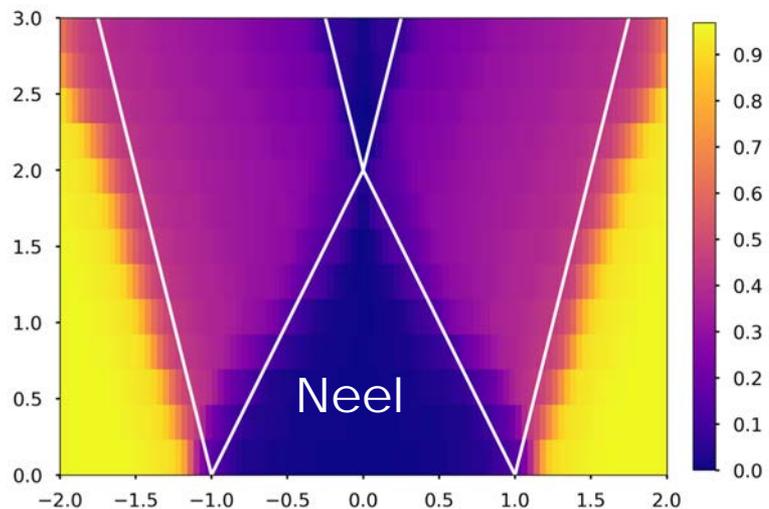
Quantum evolution Markov chain: half-cell embedding



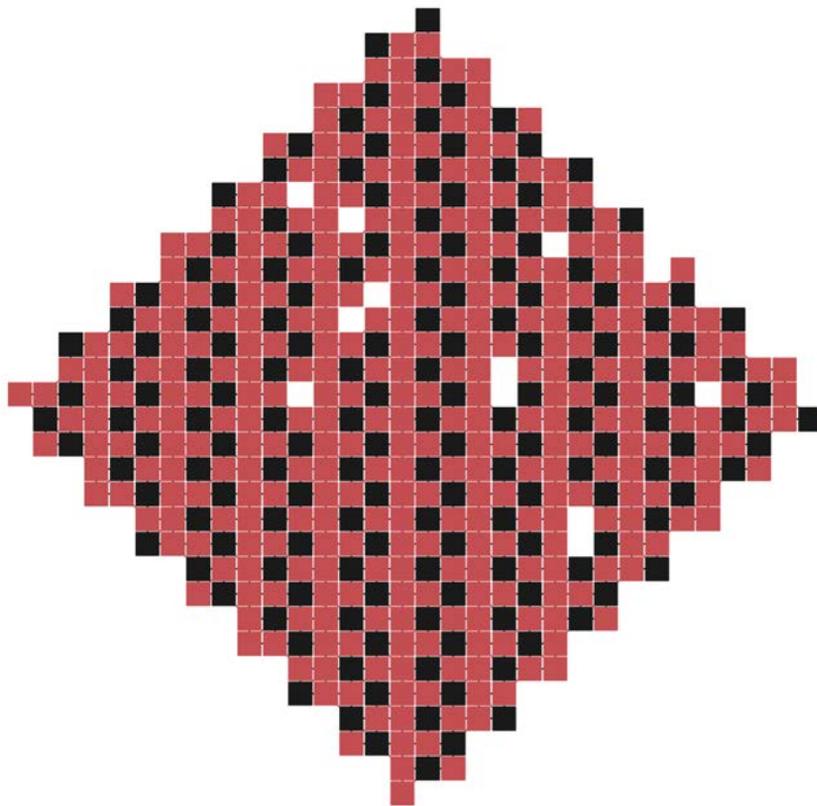
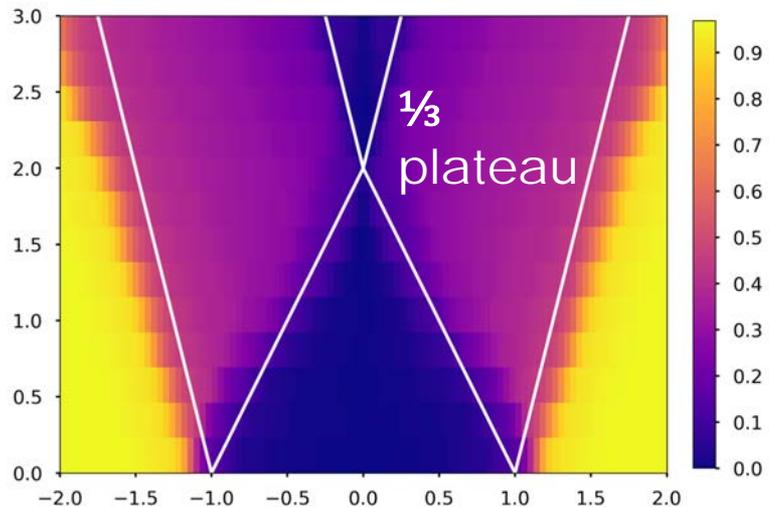
Quantum evolution Markov chain: half-cell embedding



Quantum evolution Markov chain: half-cell embedding



Quantum evolution Markov chain: half-cell embedding



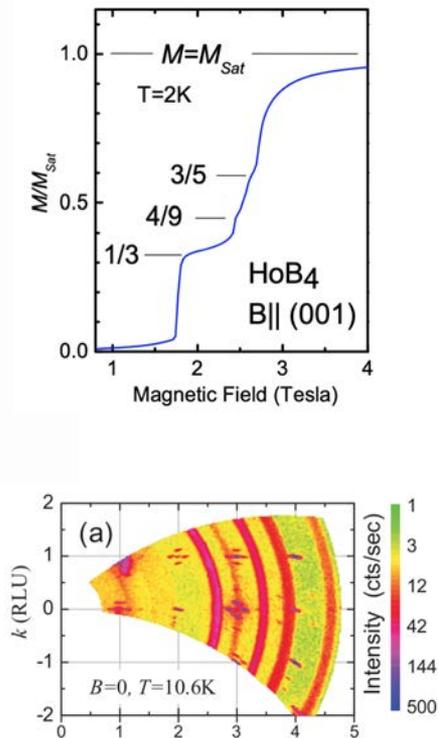
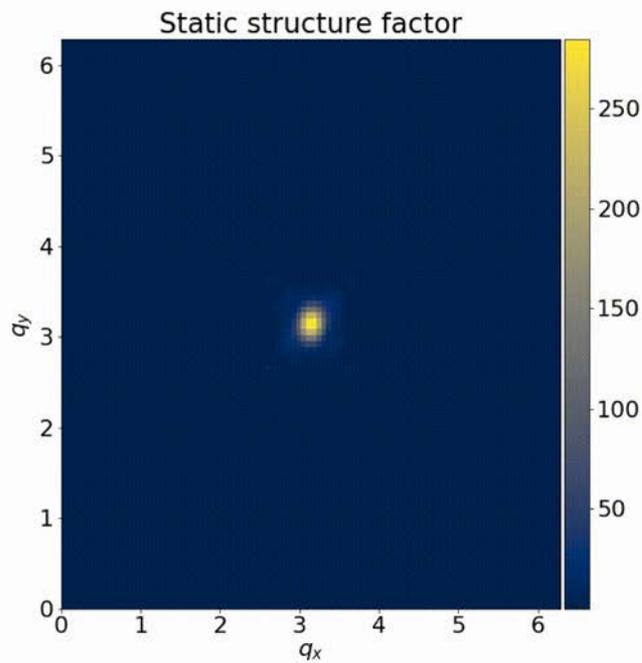
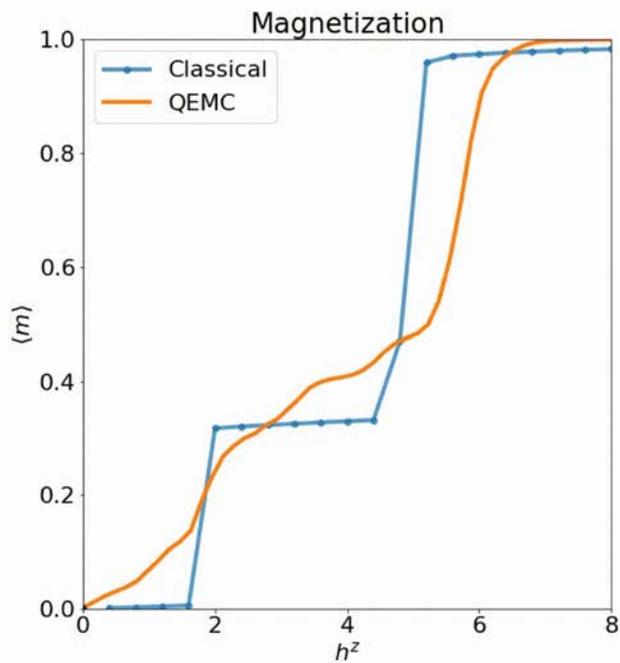
Simulations of the Ising model on a Shastry-Sutherland lattice by quantum annealing

- We demonstrate how quantum annealing enables accurate simulations of many-particle Hamiltonian systems.
- We sample the ground state energy configurations of the the SS Ising model to calculate the structure factor
- We develop a novel method for mitigating finite size and defects.
- We observe good agreement between the observed and expected material behaviors
- We can now explore defect physics and temperature effects in this model.

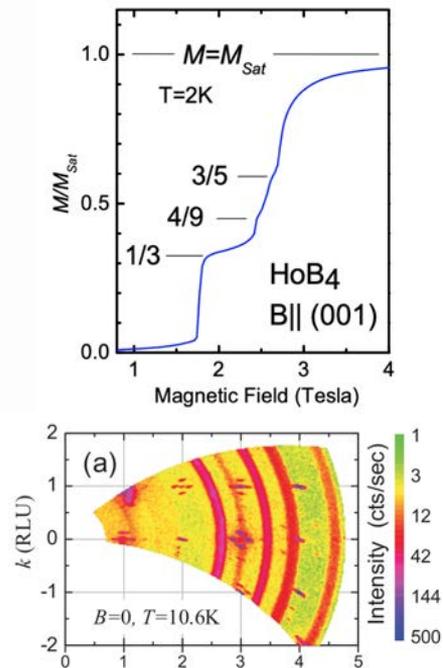
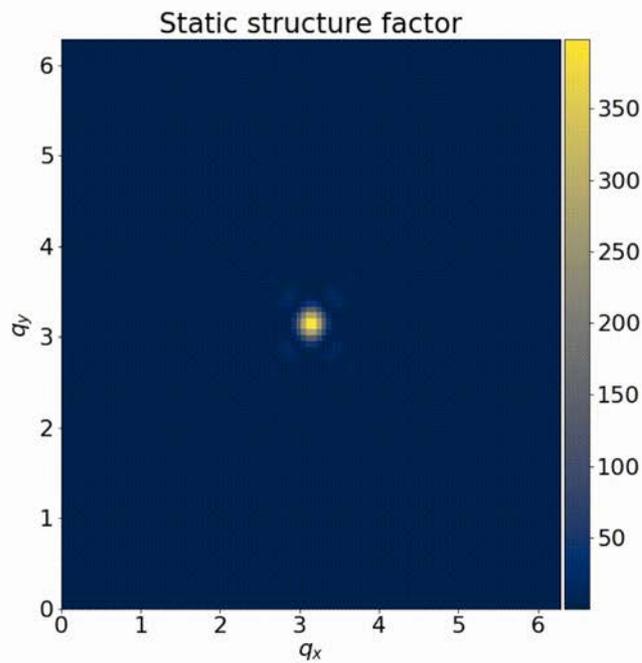
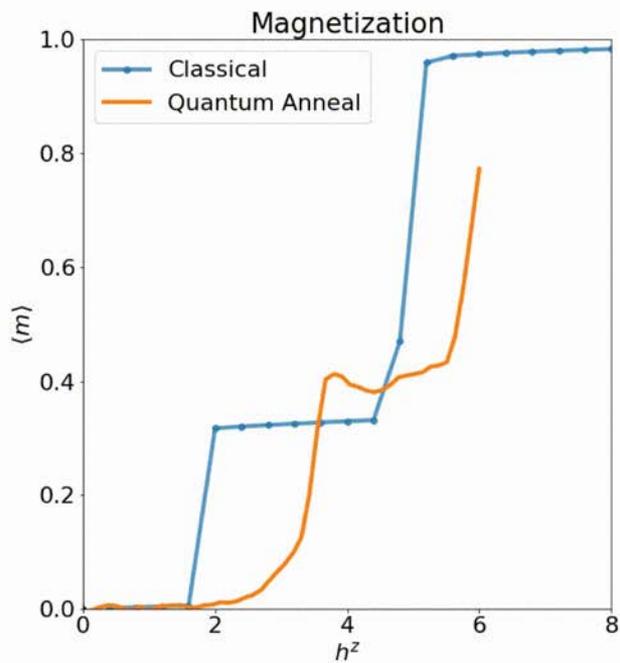
References

1. Mat’áš, S, K Siemensmeyer, E Wheeler, E Wulf, R Beyer, Th Hermannsdörfer, O Ignatchik, et al. “Magnetism of Rare Earth Tetraborides.” *Journal of Physics: Conference Series* 200, no. 3 (January 1, 2010): 032041. <https://doi.org/10.1088/1742-6596/200/3/032041>.
2. Siemensmeyer, K., E. Wulf, H.-J. Mikeska, K. Flachbart, S. Gabáni, S. Mat’áš, P. Priputen, A. Efdokimova, and N. Shitsevalova. “Fractional Magnetization Plateaus and Magnetic Order in the Shastry-Sutherland Magnet TmB₄.” *Physical Review Letters* 101, no. 17 (October 20, 2008). <https://doi.org/10.1103/PhysRevLett.101.177201>.
3. Dublenych, Yu. I. “Ground States of the Ising Model on the Shastry-Sutherland Lattice and the Origin of the Fractional Magnetization Plateaus in Rare-Earth-Metal Tetraborides.” *Physical Review Letters* 109, no. 16 (October 16, 2012). <https://doi.org/10.1103/PhysRevLett.109.167202>.

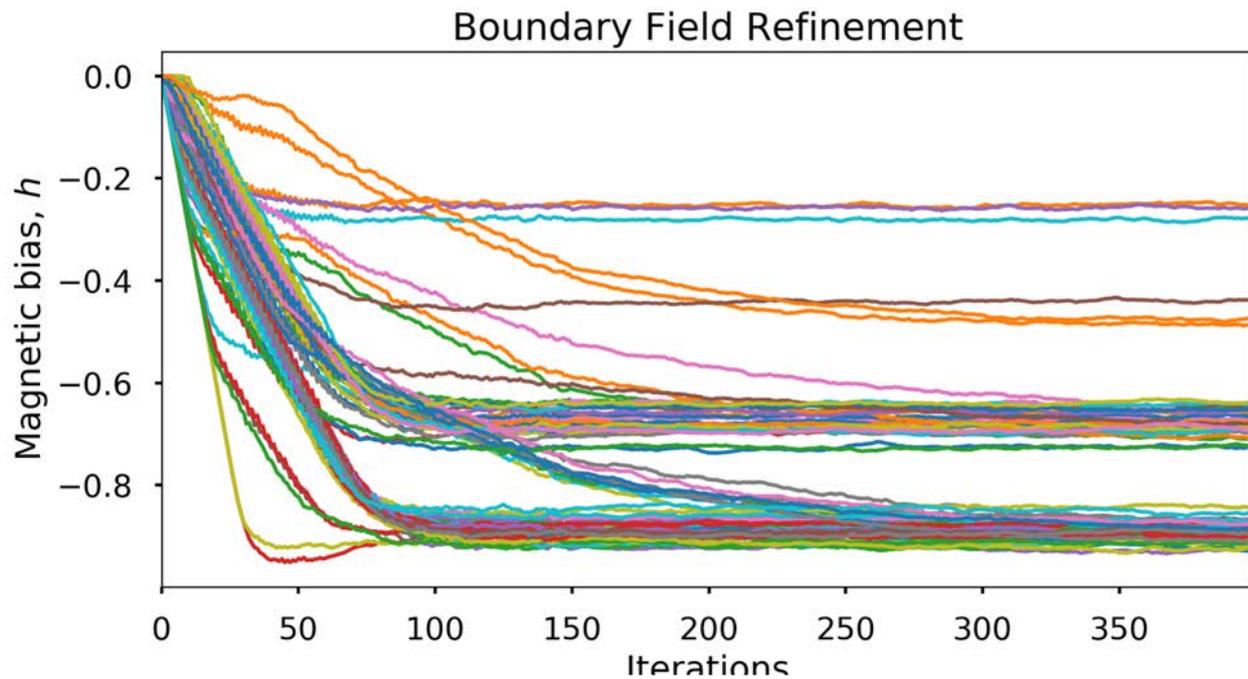
Quantum evolution Markov chain without boundary conditions



Forward Anneal - Tilt Embedding



Boundary Refinement



Quantum evolution Markov chain: half-cell embedding

